

# EE2 Mathematics : Complex Variables

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These notes are not identical word-for-word with my lectures which will be given on a BB/WB. Some of these notes may contain more examples than the corresponding lecture while in other cases the lecture may contain more detailed working. I will not be handing out copies of these notes – **you are therefore advised to attend lectures and take your own.**

1. The material in them is dependent upon the material on Complex Numbers you were taught at A-level and your 1st year.
2. Handouts are :
  - (a) **Handout No 5** of the course on “Jordan’s Lemma” in complex integration contains some of the material from §3.5.

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<sup>1</sup>Do not confuse me with Dr J. Gibbons who is also in the Mathematics Dept.

# 1 Analyticity and the Cauchy-Riemann equations

## 1.1 Derivation of the Cauchy-Riemann equations

Functions of the complex variable  $z = x + iy$

$$w = f(z) \quad (1.1)$$

are expressed in the usual manner except that the independent variable  $z = x + iy$  is complex. Thus  $f(z)$  has a real part  $u(x, y)$  and an imaginary part  $v(x, y)$

$$f(z) = u(x, y) + iv(x, y). \quad (1.2)$$

Extra difficulties appear in differentiating and integrating such functions because  $z$  varies in a plane and not on a line. For functions of a single real variable the idea of an incremental change  $\delta x$  along the  $x$ -axis has to be replaced by an incremental change  $\delta z$ . Because  $\delta z$  is a vector the question of the direction of this limit becomes an issue.

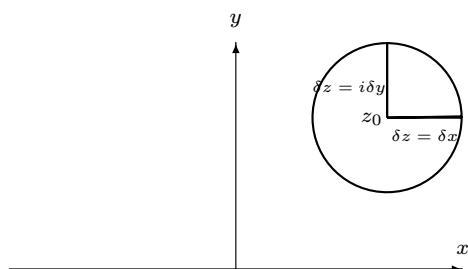
Firstly we look at the concept of differentiation. The definition of a derivative at a point  $z_0$  remains the same as usual; namely

$$\left. \frac{df(z)}{dz} \right|_{z=z_0} = \lim_{\delta z \rightarrow 0} \left( \frac{f(z_0 + \delta z) - f(z_0)}{\delta z} \right). \quad (1.3)$$

The subtlety here lies in the limit  $\delta z \rightarrow 0$  because  $\delta z$  is itself a vector and therefore the limit  $\delta z \rightarrow 0$  may be taken in many directions. If the limit in (1.3) is to be unique (to make any sense) *it is required that it be independent of the direction in which the limit  $\delta z \rightarrow 0$  is taken.* If this is the case then it is said that  $f(z)$  is differentiable at the point  $z$ .

There is a general test on functions to determine whether (1.3) is independent of the direction of the limit. The simplest way is to firstly take the limit in the horizontal direction: that is  $\delta z = \delta x$ , in which case

$$\frac{df(z)}{dz} = \frac{\partial u}{\partial x} + i \frac{\partial v}{\partial x} \equiv u_x + iv_x. \quad (1.4)$$



The  $z$ -plane with a point at  $z_0$  and a circle of radius  $|\delta z|$  around it. The horizontal radius is drawn for the case when  $\delta z = \delta x$  and the vertical for the case when  $\delta z = i\delta y$ .

Next we take the limit in the vertical direction : that is  $\delta z = i\delta y$

$$\frac{df(z)}{dz} = \frac{\partial u}{\partial(iy)} + i\frac{\partial v}{\partial(iy)} = -i\frac{\partial u}{\partial y} + \frac{\partial v}{\partial y} \equiv -iu_y + iv_y. \quad (1.5)$$

If the limits in both directions are to be equal  $df/dz$  in (1.4) and (1.5) must be equal, which makes

$$\boxed{u_x = v_y, \quad u_y = -v_x.} \quad (1.6)$$

The boxed pair of equations above are known as **the Cauchy-Riemann equations**. If these hold at a point  $z$  then  $f(z)$  is said to be differentiable at  $z$ . There is no such requirement in single variable calculus. Moreover the CR equations bring us to a further idea regarding differentiation in the complex plane :

**Definition :** *If  $f(z)$  is differentiable at all points in a neighbourhood of a point  $z_0$  then  $f(z)$  is said to be analytic (regular) at  $z_0$ .*

Some functions are analytic everywhere in the complex plane except at certain points : these points are called *singularities*. Three examples illustrate this.

**Example 1 :**  $f(z) = z^2$ . Writing  $z^2 = x^2 - y^2 + 2ixy$  we have

$$u(x, y) = x^2 - y^2, \quad v(x, y) = 2xy. \quad (1.7)$$

Clearly four trivial partial derivatives show that  $u_x = 2x$ ,  $u_y = -2y$ ,  $v_x = 2y$  and  $v_y = 2x$  thus demonstrating that the CR equations hold for all values of  $x$  and  $y$ . It follows that  $f(z) = z^2$  is differentiable at all points in the  $z$ -plane and every point in this plane has an (infinite) neighbourhood in which  $f(z) = z^2$  is differentiable. Clearly  $f(z) = z^2$  is analytic everywhere.

**Example 2 :**  $f(z) = z^{-1}$ . Writing  $z^{-1} = (x - iy)/(x^2 + y^2)$  we have

$$u(x, y) = \frac{x}{x^2 + y^2} \quad v(x, y) = -\frac{y}{x^2 + y^2}. \quad (1.8)$$

Without giving the working it is not difficult to show that the CR equations hold everywhere except at the origin  $z = 0$  where the limit is indeterminate :  $z = 0$  is the point where it fails to be differentiable. Hence  $w = z^{-1}$  is analytic everywhere except at  $z = 0$ .

**Example 3 :**  $f(z) = |z|^2$ . We have

$$u(x, y) = x^2 + y^2, \quad v(x, y) = 0, \quad (1.9)$$

and so

$$u_x = 2x, \quad u_y = 2y, \quad v_x = v_y = 0. \quad (1.10)$$

Clearly the CR equations do **not** hold anywhere except at  $z = 0$ . Therefore  $f(z) = |z|^2$  is not differentiable anywhere except at  $z = 0$  and there is no neighbourhood around  $z = 0$  in which it is differentiable. Thus the function is analytic nowhere in the  $z$ -plane.

As a final remark on this example let us look again at the limit  $\delta z \rightarrow 0$  in polar co-ordinates at a fixed point  $z_0$  about which we describe a circle  $z = z_0 + re^{i\theta}$ .

$$\begin{aligned} \left. \frac{df}{dz} \right|_{z=z_0} &= \lim_{\delta z \rightarrow 0} \left( \frac{|z + \delta z|^2 - |z|^2}{\delta z} \right)_{z=z_0} \\ &= \lim_{\delta z \rightarrow 0} \left( \frac{(\delta z)z^* + (\delta z^*)z + (\delta z)(\delta z^*)}{\delta z} \right)_{z=z_0} \\ &= z_0^* + z_0 \lim_{\delta z \rightarrow 0} \left( \frac{\delta z^*}{\delta z} \right)_{z=z_0}. \end{aligned} \quad (1.11)$$

Then (1.11) can be written as

$$\left. \frac{df}{dz} \right|_{z=z_0} = z_0^* + z_0 e^{-2i\theta}. \quad (1.12)$$

This result illustrates the problem: as  $\theta$  varies (and thus the direction of the limit) so does the limit. **This is clearly not unique except when  $z_0 = 0$ .**

## 1.2 Properties of analytic functions

Let us consider the CR equations  $u_x = v_y$  and  $u_y = -v_x$  as a condition for the analyticity of a function  $w = u(x, y) + iv(x, y)$ . Cross differentiation and elimination of first  $u$  and then  $v$  gives

$$u_{xx} + u_{yy} = 0 \quad v_{xx} + v_{yy} = 0, \quad (1.13)$$

thus showing that  $u$  and  $v$  must always be a solution of Laplace's equation (without boundary conditions): these are called **harmonic functions**. It also said that  $u(x, y)$  and  $v(x, y)$  are **conjugate** to one another. In the following set of examples it will be shown how, given a harmonic function  $u(x, y)$ , its conjugate  $v(x, y)$  can be constructed. The pair can then put together as  $u + iv = f(z)$  to ultimately find  $f(z)$ .

**Example 1:** Given that  $u = x^2 - y^2$  show (i) that it is harmonic; (ii) find  $v(x, y)$  and then (iii) construct the corresponding complex function  $f(z)$ .

With  $u = x^2 - y^2$  we have  $u_x = 2x$ ,  $u_{xx} = 2$ ,  $u_y = -2y$  and  $u_{yy} = -2$ . Therefore  $u_{xx} + u_{yy} = 0$  so it satisfies Laplace's equation. This is a sufficient condition for  $v$  to exist and for us to write  $v_y = u_x = 2x$  and  $v_x = -u_y = 2y$ . While there are two PDEs here there can only be one solution compatible with both. Integrating them both in turn gives

$$v = 2xy + A(x), \quad v = 2xy + B(y). \quad (1.14)$$

It is clear that they are compatible if  $A(x) = B(y) = \text{const} = c$  making the result

$$v = 2xy + c, \quad (1.15)$$

with

$$f(z) = x^2 - y^2 + 2ixy + ic = z^2 + ic. \quad (1.16)$$

The  $ic$  simply moves  $f(z)$  an arbitrary distance along the imaginary axis.

**Example 2:** Given that  $u = x^3 - 3xy^2$  find its conjugate function  $v(x, y)$  and the corresponding complex function  $f(z)$ .

We first check that  $u = x^3 - 3xy^2$  satisfies Laplace's equation:  $u_x = 3x^2 - 3y^2$ ;  $u_{xx} = 6x$ ;  $u_y = -6xy$  and  $u_{yy} = -6x$ . Thus  $u_{xx} + u_{yy} = 0$  and so  $v$  exists and is found from the CR equations:

$$v_y = 3x^2 - 3y^2 \quad v_x = 6xy. \quad (1.17)$$

Partially integrating these gives

$$v = 3x^2y - y^3 + A(x) \quad v = 3x^2y + B(y). \quad (1.18)$$

The way to make these compatible is to choose  $B(y) = -y^3 + c$  and  $A(x) = c$  finally giving

$$v = 3x^2y - y^3 + c \quad (1.19)$$

with

$$\begin{aligned} f(z) &= x^3 - 3xy^2 + i(3x^2y - y^3 + c) \\ &= z^3 + ic. \end{aligned} \quad (1.20)$$

**Example 3:** Given that  $u = e^x(x \cos y - y \sin y)$  show that it satisfies Laplace's equation. Also find its conjugate  $v$  and then  $f(z)$ .

We find that

$$u_{xx} = e^x[(x+2) \cos y - y \sin y]; \quad u_{yy} = -e^x[(x+2) \cos y - y \sin y], \quad (1.21)$$

and so Laplace's equation is satisfied. Then

$$v_y = u_x = e^x[(x+1) \cos y - y \sin y]; \quad v_x = -u_y = e^x[(x+1) \sin y + y \cos y]. \quad (1.22)$$

Using the indefinite integrals  $\int y \sin y \, dy = \sin y - y \cos y$  and  $\int x e^x \, dx = e^x(x-1)$  we find

$$v = e^x(x \sin y + y \cos y) + A(x); \quad v = e^x(y \cos y + x \sin y) + B(y). \quad (1.23)$$

For compatibility we take  $A(x) = B(y) = \text{const} = c$ . Then

$$\begin{aligned} w &= e^x[(x+iy) \cos y - (y-ix) \sin y] + ic \\ &= e^x[z \cos y + iz \sin y] + ic \\ &= ze^{x+iy} + ic \\ &= ze^z + ic. \end{aligned} \quad (1.24)$$

### 1.3 Orthogonality

Let us finally consider the family of curves on which  $u = \text{const}$ . From the chain rule

$$du = \frac{\partial u}{\partial x} dx + \frac{\partial u}{\partial y} dy \quad (1.25)$$

and therefore on curves of constant  $u$  we have  $du = 0$ , giving the gradient on this family as

$$\left. \frac{dy}{dx} \right|_{u=\text{const}} = -\frac{u_x}{u_y}. \quad (1.26)$$

Likewise, on the family of curves of constant  $v$

$$\left. \frac{dy}{dx} \right|_{v=\text{const}} = -\frac{v_x}{v_y} \quad (1.27)$$

giving

$$\left. \frac{dy}{dx} \right|_{u=\text{const}} \times \left. \frac{dy}{dx} \right|_{v=\text{const}} = \frac{v_x u_x}{v_y u_y}. \quad (1.28)$$

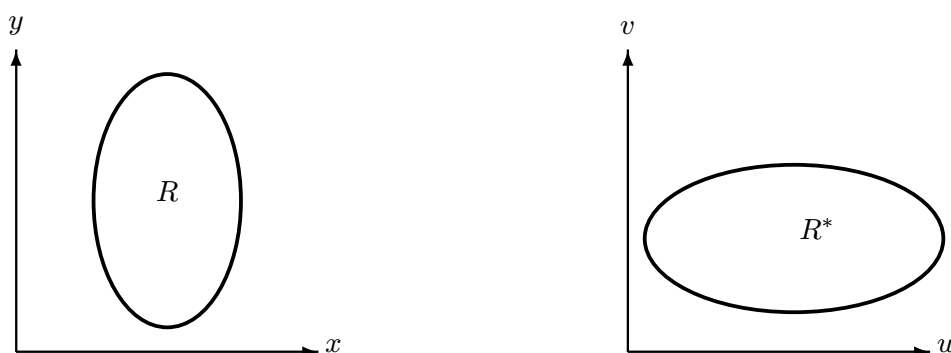
Now if  $f(z)$  is analytic in a region  $R$  then the CR equations hold there,  $u_x = v_y$  and  $u_y = -v_x$ , and (1.28) becomes

$$\left. \frac{dy}{dx} \right|_{u=\text{const}} \times \left. \frac{dy}{dx} \right|_{v=\text{const}} = -1. \quad (1.29)$$

**The final result is that in regions of analyticity curves of constant  $u$  and curves of constant  $v$  are always orthogonal.**

## 2 Mappings

### 2.1 Conformal mappings



A complex mapping  $w = f(z)$  maps a region  $R$  in the  $z$ -plane to a different region  $R^*$  in the  $w$ -plane.

A complex function  $w = f(z)$  can be thought of as a mapping from the  $z$ -plane to the  $w$ -plane. Depending on  $f(z)$  the mapping may not be unique. For instance, for  $w = z^2$  for the values  $\pm z_0$  there is one value  $w_0$ . Complex mappings do not necessarily behave in an expected

way. The concept of analyticity intrudes into these ideas in the following way. A mapping is said to be *conformal* if it preserves angles in magnitude and sense. Moreover, a mapping has a *fixed point* when  $w = f(z) = z$ . The following theorem is now stated without proof:

**Theorem 1** *The mapping defined by an analytic function  $w = f(z)$  is conformal except at points where  $f'(z) = 0$ .*

**Example 1:**  $w = z^2$  is conformal everywhere except at  $z = 0$  because  $f'(0) = 0$ . Plotting contours of  $u = x^2 - y^2$  and  $v = 2xy$  shows that conformality fails at the origin.

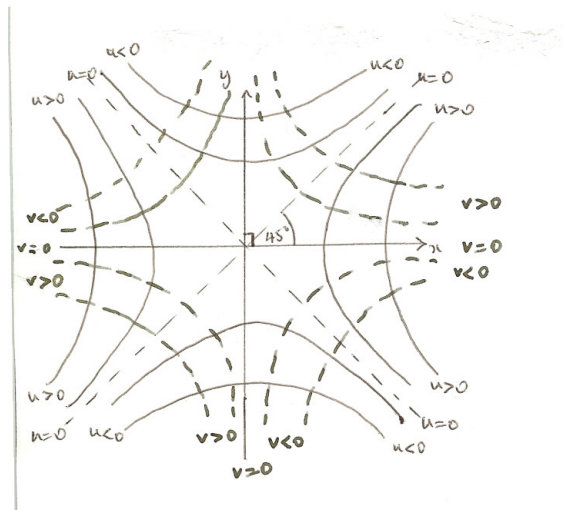


Figure 1: Contours of  $u = \text{const}$  and  $v = \text{const}$  in the  $z$ -plane: note their orthogonality except at  $z = 0$  where conformality fails.

**Example 2:** Consider  $w = \frac{1}{z-1}$  which is analytic everywhere except at  $z = 1$ . We have

$$w = \frac{1}{z-1} = \frac{1}{(x-1) + iy} = \frac{(x-1) - iy}{(x-1)^2 + y^2} \quad (2.1)$$

in which case

$$u(x, y) = \frac{x-1}{(x-1)^2 + y^2}, \quad v(x, y) = -\frac{y}{(x-1)^2 + y^2}. \quad (2.2)$$

It is clear from (2.2) that it is always true that

$$u^2 + v^2 = \frac{1}{(x-1)^2 + y^2}. \quad (2.3)$$

So far we have specified no shape in the  $z$ -plane on which this map operates. Some examples of what this map will do are these:

1. Consider the family of circles in the  $z$ -plane:  $(x - 1)^2 + y^2 = a^2$ . These circles are centred at  $(1, 0)$  of radius  $a$ . Clearly they map to

$$u^2 + v^2 = \frac{1}{a^2}, \quad (2.4)$$

which is a family of circles in the  $w$ -plane centred at  $(0, 0)$  of radius  $a^{-1}$ . As the value of  $a$  is increased the circles in the  $z$ -plane widen and those in the  $w$ -plane decrease. It is not difficult to show that the interior (exterior) of the circles in the  $z$ -plane map to the exterior (interior) of those in the  $w$ -plane. Thus we have

$$\begin{array}{l} \underline{z\text{-plane}} \quad \underline{w\text{-plane}} \\ \text{interior} \rightarrow \text{exterior} \\ \text{exterior} \rightarrow \text{interior} \end{array} \quad (2.5)$$

The circle centre  $(1, 0)$  in the  $z$ -plane maps to *the point at infinity* in the  $w$ -plane.

2. The line  $x = 0$  in the  $z$ -plane maps to what? From (2.2) and (2.3) we know that

$$u(x, y) = -\frac{1}{1 + y^2}, \quad v(x, y) = -\frac{y}{1 + y^2}, \quad u^2 + v^2 = \frac{1}{1 + y^2}. \quad (2.6)$$

$$u^2 + v^2 = \frac{1}{1 + y^2} = -u, \quad \Rightarrow \quad \left(u + \frac{1}{2}\right)^2 + v^2 = \frac{1}{4}. \quad (2.7)$$

In the  $w$ -plane this is a circle of radius  $\frac{1}{2}$  centred at  $(-\frac{1}{2}, 0)$ .

Thus we conclude that some circles can map to other circles but also straight lines can also map to circles. This is investigated in the next subsection.

## 2.2 $w = \frac{1}{z}$ maps lines/circles to lines/circles

The general equation for straight lines and circles in the  $z$ -plane can be written as

$$\alpha(x^2 + y^2) + \beta x + \gamma y + \Delta = 0. \quad (2.8)$$

where  $\alpha$ ,  $\beta$ ,  $\gamma$  and  $\Delta$  are constants. If  $\alpha = 0$  this represents a straight line but when  $\alpha \neq 0$  (2.8) represents a circle. Writing (2.8) in terms of  $z$

$$\alpha|z|^2 + \frac{\beta}{2}(z + z^*) + \frac{\gamma}{2i}(z - z^*) + \Delta = 0. \quad (2.9)$$

and then transforming to an equation  $w$  and  $w$  through  $w = \frac{1}{z}$  and  $w^* = \frac{1}{z^*}$ , (2.9) becomes

$$\frac{\alpha}{ww^*} + \frac{\beta}{2} \left( \frac{1}{w} + \frac{1}{w^*} \right) + \frac{\gamma}{2i} \left( \frac{1}{w} - \frac{1}{w^*} \right) + \Delta = 0, \quad (2.10)$$

or

$$\alpha + \frac{\beta}{2}(w + w^*) - \frac{i\gamma}{2}(w^* - w) + \Delta ww^* = 0. \quad (2.11)$$

Since  $w = u + iv$  we have

$$\alpha + \beta u - \gamma v + \Delta(u^2 + v^2) = 0. \quad (2.12)$$

This represents a family of circles in the  $u - v$  plane when  $\Delta \neq 0$  and a family of lines when  $\Delta = 0$ . Notice, that when  $\alpha \neq 0$  and  $\Delta \neq 0$  then the mapping maps circles to circles but a family of lines in the  $z$ -plane ( $\alpha = 0$ ) also maps to a family of circles in the  $w$ -plane. However, there is also the case of a family of circles in the  $z$ -plane for which  $\Delta = 0$  which map to a family of lines in the  $w$ -plane. **Thus we conclude that  $w = \frac{1}{z}$  maps lines/circles to lines/circles but not necessarily lines to lines and circles to circles.**

In addition to this we now study the *fractional linear* or *Möbius* transformation

$$w = \frac{az + b}{cz + d}, \quad ad \neq bc. \quad (2.13)$$

This includes cases such as :

(i)  $w = \frac{1}{z}$  when  $a = d = 0$ ,  $b/c = 1$ .

(ii)  $w = \frac{1}{z-1}$  as in our example above where  $a = 0$ ,  $b = 1$ ,  $c = 1$ ,  $d = -1$ .

(2.13) can be re-written as

$$w = c^{-1} \left\{ a + \frac{bc - ad}{cz + d} \right\}. \quad (2.14)$$

For various special cases :

1.  $w = z + b$ ; ( $a = d = 1$ ,  $c = 0$ ) – translation.
2.  $w = az$ ; ( $b = c = 0$ ,  $d = 1$ ) – contraction/expansion + rotation
3.  $w = \frac{1}{z}$ ; ( $a = d = 0$ ,  $b = c$ ) – maps lines/circles to lines/circles.

Thus a Möbius transformation maps lines/circles to lines/circles with contraction/expansion, rotation and translation on top.

### 2.3 Extra: Mappings of the type $w = \frac{e^z - 1}{e^z + 1}$

Consider a map  $w = \frac{e^z - 1}{e^z + 1}$  which can be re-written as

$$e^z = \frac{1 + w}{1 - w} = \frac{(1 + u + iv)(1 - u + iv)}{(1 - u)^2 + v^2}. \quad (2.15)$$

Real and imaginary parts give

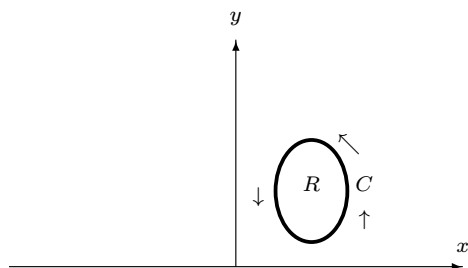
$$e^x \cos y = \frac{1 - u^2 - v^2}{(1 - u)^2 + v^2} \quad e^x \sin y = \frac{2v}{(1 - u)^2 + v^2}. \quad (2.16)$$

From these we conclude that

1. The family of lines  $y = n\pi$  in the  $z$ -plane map to the line  $v = 0$  for  $n$  integer. Thus an infinite number of horizontal lines in the  $z$ -plane all map to the  $u$ -axis in the  $w$ -plane.
2. The family of lines  $y = \frac{1}{2}(2n + 1)\pi$  in the  $z$ -plane map to the unit circle  $u^2 + v^2 = 1$  in the  $w$ -plane.

### 3 Contour Integration

#### 3.1 Cauchy's Theorem



A closed contour  $C$  enclosing a region  $R$  in the  $z$ -plane around which the line integral is considered in the counter-clockwise direction

$$\oint_C F(z) dz. \quad (3.1)$$

With

$$F(z) = u + iv \quad z = x + iy \quad (3.2)$$

we have

$$\begin{aligned} \oint_C F(z) dz &= \oint_C (u + iv)(dx + idy) \\ &= \oint_C (u dx - v dy) + i \oint_C (v dx + u dy). \end{aligned} \quad (3.3)$$

Now Green's Theorem in a plane says that for differentiable functions  $P(x, y)$  and  $Q(x, y)$

$$\oint_C (P dx + Q dy) = \int \int_R (Q_x - P_y) dx dy. \quad (3.4)$$

Therefore we have

$$\begin{aligned} \oint_C (u dx - v dy) &= \int \int_R (-v_x - u_y) dx dy, \\ \oint_C (v dx + u dy) &= \int \int_R (u_x - v_y) dx dy, \end{aligned} \quad (3.5)$$

which turns (3.3) into

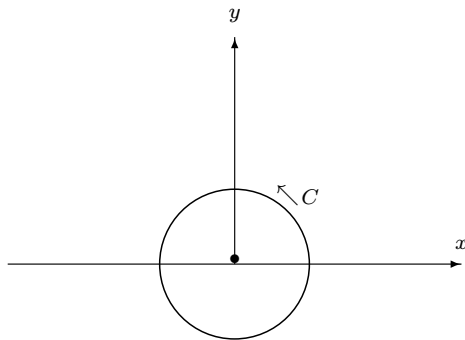
$$\oint_C F(z) dz = - \int \int_R (v_x + u_y) dx dy + i \int \int_R (u_x - v_y) dx dy. \quad (3.6)$$

**If  $F(z)$  is analytic everywhere within and on  $C$  then  $u$  and  $v$  must satisfy the CR equations:**  $u_x = v_y$  and  $v_x = -u_y$ , in which case both the real and imaginary parts on the RHS of (3.6) must be zero. We have established *Cauchy's Theorem*:

**Theorem 2** *If  $F(z)$  is analytic everywhere within and on a closed, piecewise smooth contour  $C$  then*

$$\oint_C F(z) dz = 0. \quad (3.7)$$

The key point is that provided  $F(z)$  is analytic everywhere within and on  $C$  singularities in  $F(z)$  *outside* of  $C$  are irrelevant.



For  $F(z) = z^{-1}$ , the circular contour  $C$  of radius  $a$  encloses a singularity  $\bullet$  at the origin in the  $z$ -plane. The line integral is no longer zero because of this singularity.

Now write the circular contour  $C$  as  $z = a \exp(i\theta)$  for  $\theta : 0 \rightarrow 2\pi$

$$\oint_C \frac{dz}{z} = \int_0^{2\pi} \frac{ia \exp(i\theta) d\theta}{a \exp(i\theta)} = i \int_0^{2\pi} d\theta = 2\pi i. \quad (3.8)$$

The singularity at  $z = 0$  contributes a non-zero value of  $2\pi i$  to the integral. Note that it is independent of the value of  $a$ , which is consistent with this being the only non-zero contribution to the integral.

Given the powerful result of Cauchy's Theorem, our task is to see, in a more formal manner, how singularities contribute to complex integrals. Before this their nature and classification is necessary.

### 3.2 Poles and Residues

Singularities of complex functions can take many forms but the simplest class is what are called **simple poles**. A function  $F(z)$  has a simple pole at  $z = a$  (which could be real) if it can be written in the form

$$F(z) = \frac{f(z)}{z - a} \quad (3.9)$$

where  $f(z)$  is an analytic function.  $F(z)$  has a pole of multiplicity  $m$  at  $z = a$  if it can be written in the form

$$F(z) = \frac{f(z)}{(z - a)^m} \quad (3.10)$$

where  $m = 1, 2, 3, 4, \dots$ : when  $m = 2$  we have a double pole. While poles are singularities not all singularities are poles. For instance,  $\ln z$  has a singularity at  $z = 0$  but this is not a pole nor is it a pole when  $m$  is not an integer.

**Definition 1:** The **residue** of  $F(z)$  at a simple pole at  $z = a$  is

$$\text{Residue of } F(z) = \lim_{z \rightarrow a} \{(z - a)F(z)\}. \quad (3.11)$$

**Definition 2:** The **residue** of  $F(z)$  at a pole of multiplicity<sup>2</sup>  $m$  at  $z = a$  is

$$\text{Residue of } F(z) = \lim_{z \rightarrow a} \left\{ \frac{1}{(m-1)!} \frac{d^{m-1}}{dz^{m-1}} [(z-a)^m F(z)] \right\}. \quad (3.12)$$

Note that a function may have many poles and each pole has its own residue.

**Example 1:** Consider  $F(z) = \frac{2z}{(z-1)(z-2)}$  which has simple poles at  $z = 1$  and  $z = 2$

$$\text{Residue at } z = 1 = \lim_{z \rightarrow 1} \{(z-1)F(z)\} = -2. \quad (3.13)$$

$$\text{Residue at } z = 2 = \lim_{z \rightarrow 2} \{(z-2)F(z)\} = 4. \quad (3.14)$$

**Example 2:**  $F(z) = \frac{2z}{(z-1)^2(z+4)}$  has a double pole at  $z = 1$  and a simple pole at  $z = -4$ .

$$\text{Residue at } z = -4 = \lim_{z \rightarrow -4} \{(z+4)F(z)\} = -8/25. \quad (3.15)$$

$$\begin{aligned} \text{Residue at double pole at } z = 1 &= \lim_{z \rightarrow 1} \frac{1}{1!} \left\{ \frac{d}{dz} [(z-1)^2 F(z)] \right\} \\ &= \lim_{z \rightarrow 1} \frac{d}{dz} \left[ \frac{2z}{z+4} \right] \\ &= 2 \lim_{z \rightarrow 1} \left[ \frac{(z+4) - z}{(z+4)^2} \right] = 8/25. \end{aligned} \quad (3.16)$$

It is of no significance that the residues have opposite signs.

**3.3** The residue of  $F(z) = \frac{h(z)}{g(z)}$  when  $g(z)$  has a simple zero at  $z = a$

We expand  $g(z)$  about its zero at  $z = a$  in a Taylor series

$$g(z) = g(a) + (z-a)g'(a) + \frac{1}{2}(z-a)^2 g''(a) + \dots \quad (3.17)$$

Thus, noting that  $g(a) = 0$  we have

$$\begin{aligned} \text{Residue at } z = a &= \lim_{z \rightarrow a} \left\{ \frac{(z-a)h(z)}{g(z)} \right\} \\ &= \lim_{z \rightarrow a} \left\{ \frac{(z-a)h(z)}{(z-a)g'(a) + \frac{1}{2}(z-a)^2 g''(a) + \dots} \right\} \\ &= \lim_{z \rightarrow a} \left\{ \frac{h(z)}{g'(a) + \frac{1}{2}(z-a)g''(a) + \dots} \right\} = \frac{h(a)}{g'(a)}. \end{aligned} \quad (3.18)$$

---

<sup>2</sup>This formula will be quoted at the bottom of an exam question: it is found from a coefficient in what is known as a Laurent expansion – see Kreysig's book.

### 3.4 The Residue Theorem

Now consider a simple pole at  $z = a$  as in the Figure below. As explained in the caption, the pole is isolated by a device which consists of taking a “cut” into the contour and inscribing a small circle of radius  $r$  around it. The full *closed* contour  $C$  consists of the two edges of the cuts  $C^\pm$  running in opposite directions, the small circle  $C_a$  and then  $C_1$  which is the rest of  $C$  with the small piece  $\epsilon$  removed. Thus we have

$$C : C_1 + C_a + C^+ + C^- . \quad (3.19)$$

This device ensures that the pole lies outside of  $C$  as it has been constructed, in which case Cauchy’s Theorem says that

$$\oint_C F(z) dz = 0 \quad (3.20)$$

in which case

$$0 = \int_C F(z) dz = \left( \int_{C_1} + \int_{C_a} + \int_{C^+} + \int_{C^-} \right) F(z) dz . \quad (3.21)$$

Two points to note are :

1. The four integrals are not closed so they don’t have the  $\oint$  notation. In the limit  $\epsilon \rightarrow 0$  the integrals  $\int_{C^+}$  and  $\int_{C^-}$  cancel as they go in opposite directions ;
2. The integral over  $C_a$  is clockwise, not counter-clockwise.

We are left with

$$\lim_{\epsilon \rightarrow 0} \int_{C_1} F(z) dz = - \lim_{\epsilon \rightarrow 0} \int_{C_a \leftarrow} F(z) dz = \int_{C_a \rightarrow} F(z) dz \quad (3.22)$$

We now write

$$\int_{C_a \rightarrow} F(z) dz = \int_{C_a \rightarrow} \frac{f(z)}{z - a} dz . \quad (3.23)$$

As in the figure we write the equation of the circle of radius  $a$  in the complex plane as  $z = a + re^{i\theta}$  so

$$\int_{C_a \rightarrow} F(z) dz = \int_{C_a \rightarrow} \frac{f(z)}{z - a} dz = \int_{C_a \rightarrow} \frac{f(a + re^{i\theta})}{re^{i\theta}} ire^{i\theta} d\theta . \quad (3.24)$$

Our next step at this stage is to take the limit  $r \rightarrow 0$

$$\lim_{r \rightarrow 0} \int_{C_a \rightarrow} F(z) dz = \lim_{r \rightarrow 0} \int_{C_a \rightarrow} \frac{f(a + re^{i\theta})}{re^{i\theta}} ire^{i\theta} d\theta = 2\pi i f(a) \quad (3.25)$$

which gives **Cauchy’s integral formula**

$$\oint_C F(z) dz = \lim_{r \rightarrow 0} \int_{C_a \rightarrow} F(z) dz = 2\pi i f(a) . \quad (3.26)$$

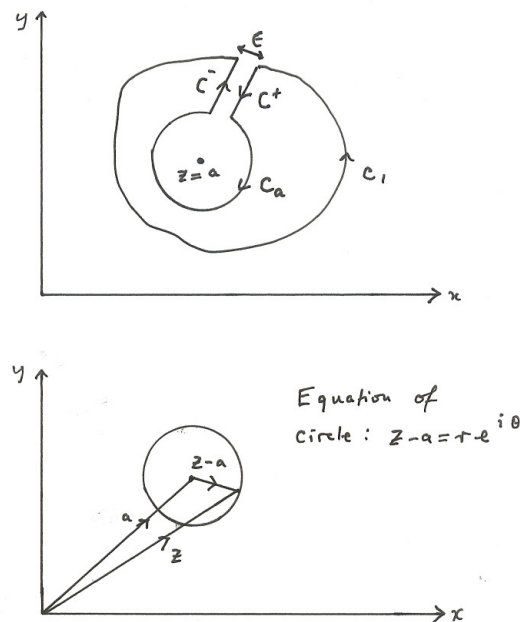


Figure 2: The top figure is the full contour  $C$  comprising the two edges of the cuts  $C^\pm$  running in opposite directions, the small circle  $C_a$  and then  $C_1$  which is the rest of  $C$  with the small piece  $\epsilon$  removed. The bottom figure shows why the equation of the small circle of radius  $r$  is  $z = a + r \exp i\theta$ .

However, because  $z = a$  is a simple pole

$$\text{At the pole at } z = a \text{ the residue of } F(z) = \lim_{z \rightarrow a} \{(z - a)F(z)\} = f(a) \quad (3.27)$$

We have proved

$$\oint_C F(z) dz = 2\pi i \times \{\text{Residue of } F(z) \text{ at the simple pole at } z = a\}. \quad (3.28)$$

It is clear that this procedure of making a cut and ring-fencing a pole can be performed for many simple poles and the individual residues added. The result can also be proved (not here) when poles have higher multiplicity. Altogether we have proved :

**Theorem 3 (Residue Theorem :)** *If the only singularities of  $F(z)$  within  $C$  are poles then*

$$\oint_C F(z) dz = 2\pi i \times \{\text{Sum of the residues of } F(z) \text{ at its poles within } C\}. \quad (3.29)$$

Some examples of this immensely powerful theorem follow.

**Example 1:** Find

$$\oint_{C_i} \frac{2z dz}{(z-1)(z-2)} \quad (3.30)$$

where (i)  $C_1$  is the circle centred at  $(0, 0)$  of radius 3 and (ii)  $C_2$  is the circle centred at  $(0, 0)$  of radius  $3/2$ .

$F(z)$  has two simple poles: the first at  $z = 1$  and the second at  $z = 2$ . Their residues have been found in (3.13) and (3.14). For  $C_1$  both poles lie inside  $C_1$

$$\oint_{C_1} F(z) dz = 2\pi i \times (-2 + 4) = 4\pi i, \quad (3.31)$$

whereas for  $C_2$  only the pole at  $z = 1$  lie inside, thus

$$\oint_{C_2} F(z) dz = 2\pi i \times (-2) = -4\pi i. \quad (3.32)$$

**Example 2:** Find

$$\oint_C \frac{2z dz}{(z-1)^2(z+4)} \quad (3.33)$$

where  $C$  is the circle of radius 5 centred at  $z = 0$ . For this  $C$  both poles lie inside so both must be included. From the residues computed in (3.15) and (3.16) we find that

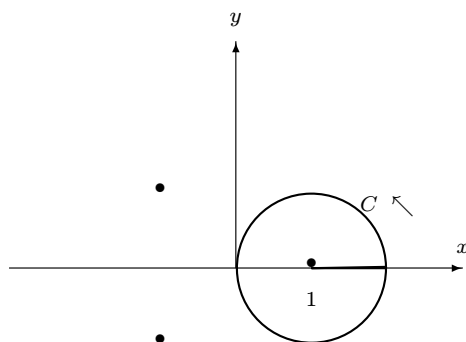
$$\oint_C \frac{2z dz}{(z-1)^2(z+4)} = 2\pi i \times \{-8/25 + 8/25\} = 0. \quad (3.34)$$

As remarked before, the zero sum of the residues has no significance.

**Example 3:** Find

$$\oint_C \frac{dz}{(z^3 - 1)^2} \quad (3.35)$$

where  $C$  is the circle  $|z - 1| = 1$ .



The contour is the circle  $|z - 1| = 1$  in the  $z$ -plane. The double pole lies at  $z = 1$  whereas the two other double poles lie outside  $C$  at  $z = \exp 2\pi i/3$  and  $z = \exp -2\pi i/3$ .

$z^3 - 1 = 0$  factors into  $(z - 1)(z^2 + z + 1) = 0$  so it has zeroes at 1,  $z = \exp(\pm 2\pi i/3)$ . These are double poles for  $F(z)$  but only the double pole at  $z = 1$  lies inside  $C$ . Its residue there is

$$\begin{aligned} \lim_{z \rightarrow 1} \frac{d}{dz} \left\{ \frac{(z-1)^2}{(z^3-1)^2} \right\} &= \lim_{z \rightarrow 1} \frac{d}{dz} \left\{ \frac{1}{(z^2+z+1)^2} \right\} \\ &= -2 \lim_{z \rightarrow 1} \left\{ \frac{2z+1}{(z^2+z+1)^3} \right\} = -2/9. \end{aligned} \quad (3.36)$$

Therefore we deduce from the Residue Theorem that

$$\oint_C \frac{dz}{(z^3-1)^2} = -4\pi i/9. \quad (3.37)$$

**Example 4: (2006)** Find

$$\oint_C \frac{z dz}{(z-1)^2(z-i)} \quad (3.38)$$

where  $C$  is the circle  $|z| = 2$ .

For the double pole at  $z = 1$ , the residue there is

$$\begin{aligned} \lim_{z \rightarrow 1} \frac{d}{dz} \left\{ \frac{z(z-1)^2}{(z-1)^2(z-i)} \right\} &= \lim_{z \rightarrow 1} \frac{d}{dz} \left\{ \frac{z}{z-i} \right\} \\ &= -\frac{i}{(1-i)^2} = \frac{1}{2}. \end{aligned} \quad (3.39)$$

For the simple pole at  $z = i$  the residue there is

$$\lim_{z \rightarrow i} \left\{ \frac{z(z-i)}{(z-1)^2(z-i)} \right\} = \frac{i}{(1-i)^2} = -\frac{1}{2}. \quad (3.40)$$

Both poles must be included within  $C$  so we conclude from the Residue Theorem that

$$\oint_C \frac{z dz}{(z-1)^2(z-i)} = \left(\frac{1}{2} - \frac{1}{2}\right) = 0. \quad (3.41)$$

**Example 5: (2006)** Find

$$\oint_C \frac{z^2 dz}{(z-i)^3} \quad (3.42)$$

where  $C$  is the circle  $|z| = 2$  as above.

For the triple pole at  $z = i$  the residue there is

$$\lim_{z \rightarrow i} \frac{1}{2!} \frac{d^2}{dz^2} \left\{ \frac{z^2(z-i)^3}{(z-i)^3} \right\} = 1. \quad (3.43)$$

Hence

$$\oint_C \frac{z^2 dz}{(z-i)^3} = 2\pi i. \quad (3.44)$$

### 3.5 Improper integrals and Jordan's Lemma

The Residue Theorem can be used to evaluate *real integrals* of the type

$$\int_{-\infty}^{\infty} e^{imx} F(x) dx, \quad m \geq 0, \quad (3.45)$$

provided  $F(x)$  has certain convergence properties: these are called *improper integrals* because of the infinite nature of their limits. Formally we write them as

$$\int_{-\infty}^{\infty} e^{imx} F(x) dx = \lim_{R \rightarrow \infty} \int_{-R}^R e^{imx} F(x) dx. \quad (3.46)$$

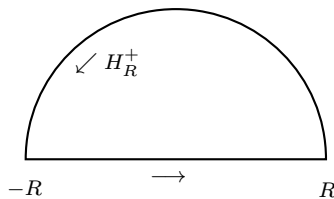
The main idea is to consider a class of *complex integrals*

$$\oint_C e^{imz} F(z) dz \quad (3.47)$$

where  $C$  consists of the semi-circular arc as in the figure below. The two essential parts are the arc of the semicircle of radius  $R$  denoted by  $H_R$  and the real axis  $[-R, R]$ . Hence we can re-write (3.47) as

$$\oint_C e^{imz} F(z) dz = \underbrace{\int_{-R}^R e^{imx} F(x) dx}_{\text{real integral}} + \underbrace{\int_{H_R} e^{imz} F(z) dz}_{\text{complex integral}} \quad (3.48)$$

In principle the closed complex integral over  $C$  on the LHS can be evaluated by the Residue Theorem. Our next aim is to evaluate the real integral on the RHS in the limit  $R \rightarrow \infty$ . This requires a result which is called Jordan's Lemma.



Jordan's Lemma deals with the problem of how a contour integral behaves on the semi-circular arc  $H_R^+$  of a closed contour  $C$ .

**Lemma 1 (Jordan)** *If the only singularities of  $F(z)$  are poles, then*

$$\lim_{R \rightarrow \infty} \int_{H_R} e^{imz} F(z) dz = 0 \quad (3.49)$$

*provided that  $m > 0$  and  $|F(z)| \rightarrow 0$  as  $R \rightarrow \infty$ . If  $m = 0$  then a faster convergence to zero is required for  $F(z)$ .*

**Proof:** Since  $H_R$  is the semi-circle  $z = Re^{i\theta} = R(\cos \theta + i \sin \theta)$  and  $dz = iRe^{i\theta} d\theta$

$$\begin{aligned} \lim_{R \rightarrow \infty} \left| \int_{H_R} e^{imz} F(z) dz \right| &= \lim_{R \rightarrow \infty} \left| \int_{H_R} e^{imR \cos \theta - mR \sin \theta} F(z) R e^{i\theta} d\theta \right| \\ &\leq \lim_{R \rightarrow \infty} \int_{H_R} e^{-mR \sin \theta} |F(z)| R d\theta \end{aligned} \quad (3.50)$$

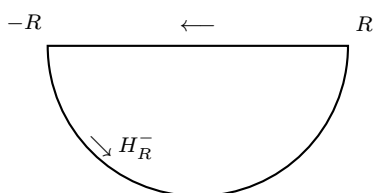
having recalled that  $|e^{i\alpha}| = 1$  for any real  $\alpha$  and  $|\int f(z) dz| \leq \int |f(z)| dz$ . Note that in the exp-term on the RHS of (3.50),  $\sin \theta > 0$  in the upper half plane. Hence, provided  $m > 0$ , the exponential ensures that the RHS is zero in the limit  $R \rightarrow \infty$  (see remarks below).  $\square$

### Remarks:

a) When  $m > 0$  forms of  $F(z)$  such as  $F(z) = \frac{1}{z}$ ,  $F(z) = \frac{1}{z+a}$  or rational functions of  $z$  such as  $F(z) = \frac{z^p \dots}{z^q + \dots}$  (for  $0 \leq p < q$  and  $p$  and  $q$  integers) will all converge fast enough as these all have simple poles and  $|F(z)| \rightarrow 0$  as  $R \rightarrow \infty$ .

b) If, however,  $m = 0$  then a modification is needed: e.g. if  $F(z) = \frac{1}{z}$  then  $|F(z)| \rightarrow 0$  but  $\lim_{R \rightarrow \infty} z|F(z)| = 1$ . We need to alter the restriction on the integers  $p$  and  $q$  to  $0 \leq p < q - 1$  which excludes cases like  $F(z) = \frac{1}{z}$ ,  $F(z) = \frac{1}{z+a}$ .

c) What about  $m < 0$ ? To ensure that the exponential is decreasing for  $R \rightarrow \infty$  we need  $\sin \theta < 0$ . This is true in the lower half plane. Hence in this case we take our contour in the lower half plane (call this  $H_R^-$  as opposed to  $H_R^+$  in the upper) but still in an anti-clockwise direction.



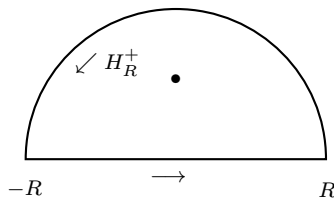
A contour in the lower  $\frac{1}{2}$ -plane with semi-circle  $H_R^-$  taken in the counter-clockwise direction which is used for cases when  $m < 0$ . See the notes on Fourier Transforms for cases when this is useful.

The conclusion is that if  $F(z)$  satisfies the conditions for Jordan's Lemma then

$$\int_{-\infty}^{\infty} e^{imx} F(x) dx = 2\pi i \times \{\text{Sum of the residues of the poles of } e^{imz} F(z) \text{ in the upper } \frac{1}{2}\text{-plane}\}. \quad (3.51)$$

**Example 1:** Show that

$$\int_{-\infty}^{\infty} \frac{dx}{1+x^2} = \pi. \quad (3.52)$$



$C$  is comprised of a semi-circular arc  $H_R^+$  and a section on the real axis from  $-R$  to  $R$ . Only the simple pole at  $z = i$  lies within  $C$ .

Thus we consider the complex integral over the semicircle  $C$  in the upper half-plane

$$\oint_C \frac{dz}{1+z^2} \quad (3.53)$$

with  $m = 0$ . The simple pole at  $z = i$  and the quadratic nature of the denominator is enough for convergence and so from Jordan's Lemma

$$\lim_{R \rightarrow \infty} \int_{H_R} \frac{dz}{1+z^2} = 0. \quad (3.54)$$

The residue of  $F(z)$  at the pole in the upper-half-plane at  $z = i$  is  $1/2i$  and so from the Residue Theorem<sup>3</sup>

$$\oint_C \frac{dz}{1+z^2} = 2\pi i \times 1/2i = \pi. \quad (3.55)$$

Finally we have the result

$$\pi = \lim_{R \rightarrow \infty} \int_{-R}^R \frac{dx}{1+x^2}. \quad (3.56)$$

**Example 2:** Show that

$$\int_{-\infty}^{\infty} \frac{dx}{1+x^4} = \pi/\sqrt{2}. \quad (3.57)$$

We consider the complex integral over the semicircle  $C$  in the upper half-plane

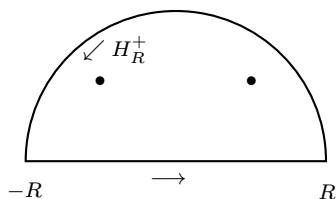
$$\oint_C \frac{dz}{1+z^4} \quad (3.58)$$

with  $m = 0$ . The existence of poles as the only singularities and the quartic nature of the denominator allows us to appeal to Jordan's Lemma

$$\lim_{R \rightarrow \infty} \int_{H_R} \frac{dz}{1+z^4} = 0. \quad (3.59)$$

---

<sup>3</sup>Note that this result could have been found by direct integration but this can only be done for the case  $n = 1$  when the denominator is  $1 + x^{2n}$ . See Examples 2 and 3.



The only simple poles in the upper half-plane at  $e^{i\pi/4}$ ,  $e^{3i\pi/4}$  lie within  $C$ .

$z^4 = -1$  has four zeroes lying at  $e^{i\pi/4}$ ,  $e^{3i\pi/4}$  in the upper half-plane and  $e^{-i\pi/4}$ ,  $e^{-3i\pi/4}$  in the lower half-plane. Only the first two are relevant. Now we use the trick in (3.18) to find the residues of the two poles in the upper half-plane. Using  $h(z) = 1$  and  $g(z) = 1 + z^4$  the residues at  $e^{i\pi/4}$  and  $e^{3i\pi/4}$  are

$$\frac{1}{4z^3} \Big|_{z=e^{i\pi/4}} = \frac{1}{4}e^{-3i\pi/4} \quad \text{and} \quad \frac{1}{4z^3} \Big|_{z=e^{3i\pi/4}} = \frac{1}{4}e^{-9i\pi/4} \quad (3.60)$$

Thus our final result is

$$\begin{aligned} \int_{-\infty}^{\infty} \frac{dx}{1+x^4} &= \frac{1}{2}\pi i (e^{-3i\pi/4} + e^{-9i\pi/4}) \\ &= \frac{1}{2}\pi i (-e^{i\pi/4} + e^{-i\pi/4}) \\ &= \pi \sin\left(\frac{\pi}{4}\right) = \pi/\sqrt{2}. \end{aligned} \quad (3.61)$$

**Example 3:** Show that

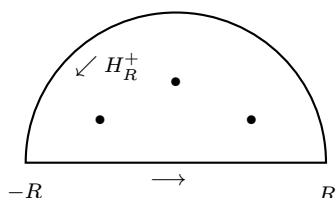
$$\int_{-\infty}^{\infty} \frac{dx}{1+x^6} = 2\pi/3. \quad (3.62)$$

Thus we consider the complex integral over the semicircle  $C$  in the upper half-plane

$$\oint_C \frac{dz}{1+z^6} \quad (3.63)$$

with  $m = 0$ . The sextic nature of the denominator is enough for fast convergence and so from Jordan's Lemma

$$\lim_{R \rightarrow \infty} \int_{H_R} \frac{dz}{1+z^6} = 0. \quad (3.64)$$



The only simple poles in the upper half-plane at  $e^{i\pi/6}$ ,  $e^{i\pi/2}$  and  $e^{5i\pi/6}$  lie within  $C$ .

$z^6 = -1$  has six zeroes lying at  $e^{i\pi/6}$ ,  $e^{i\pi/2}$  and  $e^{5i\pi/6}$  in the upper half-plane and a further three in the lower half-plane. Now we use the trick in (3.18) to find the residues of the three

poles in the upper half-plane. Using  $h(z) = 1$  and  $g(z) = 1 + z^6$  the residues at  $e^{i\pi/6}$  and  $e^{i\pi/2}$  and  $e^{5i\pi/6}$  are

$$\frac{1}{6z^5} \Big|_{z=e^{i\pi/6}} = \frac{1}{6} e^{-5i\pi/6}; \quad \frac{1}{6z^5} \Big|_{z=e^{i\pi/2}} = \frac{1}{6} e^{-5i\pi/2}; \quad \frac{1}{6z^5} \Big|_{z=e^{5i\pi/6}} = \frac{1}{6} e^{-25i\pi/6} \quad (3.65)$$

Thus our final result is

$$\begin{aligned} \int_{-\infty}^{\infty} \frac{dx}{1+x^6} &= \frac{2\pi i}{6} (e^{-5i\pi/6} + e^{-5i\pi/2} + e^{-i\pi/6}) \\ &= -\frac{2\pi i^2}{6} (2 \sin \frac{1}{6}\pi + \sin \pi/2) = 2\pi/3. \end{aligned} \quad (3.66)$$

**Example 4:** For  $m > 0$  show that

$$\int_{-\infty}^{\infty} \frac{\cos mx \, dx}{a^2 + x^2} = \frac{\pi}{a} e^{-ma}. \quad (3.67)$$

We consider the complex integral over the semicircle  $C$  in the upper half-plane

$$\oint_C \frac{e^{imz} dz}{a^2 + z^2}. \quad (3.68)$$

The integrand has only one simple pole at  $z = ia$  in the upper half-plane whose residue is

$$\lim_{z \rightarrow ia} \left\{ \left( \frac{z - ia}{a^2 + z^2} \right) e^{imz} \right\} = \frac{e^{-ma}}{2ia}. \quad (3.69)$$

Therefore, from the Residue Theorem

$$\oint_C \frac{e^{imz} dz}{a^2 + z^2} = 2\pi i \times \frac{e^{-ma}}{2ia} = \frac{\pi}{a} e^{-ma}. \quad (3.70)$$

Moreover, from Jordan's Lemma

$$\lim_{R \rightarrow \infty} \int_{H_R} \frac{e^{imz} dz}{a^2 + z^2} = 0. \quad (3.71)$$

Therefore

$$\frac{\pi}{a} e^{-ma} = \int_{-\infty}^{\infty} \frac{e^{imx} dx}{a^2 + x^2} + 0. \quad (3.72)$$

What happens to the imaginary part

$$\int_{-\infty}^{\infty} \frac{\sin mx \, dx}{a^2 + x^2} ? \quad (3.73)$$

Notice that the integrand is an *odd function*: therefore, over  $(-\infty, \infty)$  the part over  $(-\infty, 0)$  will cancel with that over  $(0, \infty)$ , leaving zero as a result. Thus we have

$$\int_{-\infty}^{\infty} \frac{\cos mx \, dx}{a^2 + x^2} = \frac{\pi}{a} e^{-ma}. \quad (3.74)$$

**Example 5:** For  $m > 0$  show that

$$\int_{-\infty}^{\infty} \frac{\cos mx \, dx}{(a^2 + x^2)^2} = \frac{\pi e^{-ma}}{2a^3} (1 + ma). \quad (3.75)$$

We consider the complex integral over the semicircle  $C$  in the upper half-plane

$$\oint_C \frac{e^{imz} dz}{(a^2 + z^2)^2}. \quad (3.76)$$

The integrand has only one double pole at  $z = ia$  in the upper half-plane whose residue is

$$\lim_{z \rightarrow ia} \frac{d}{dz} \left[ \frac{(z - ia)^2 e^{imz}}{(a^2 + z^2)^2} \right] = -\frac{ie^{-ma}}{4a^3} (1 + ma). \quad (3.77)$$

Therefore, from the Residue Theorem

$$\begin{aligned} \oint_C \frac{e^{imz} dz}{(a^2 + z^2)^2} &= -2\pi i \times \frac{ie^{-ma}}{4a^3} (1 + ma) \\ &= \frac{\pi e^{-ma}}{2a^3} (1 + ma) \end{aligned} \quad (3.78)$$

Moreover, from Jordan's Lemma, with  $m > 0$

$$\lim_{R \rightarrow \infty} \int_{H_R} \frac{e^{imz} dz}{(a^2 + z^2)^2} = 0. \quad (3.79)$$

Therefore

$$\frac{\pi e^{-ma}}{2a^3} (1 + ma) = \int_{-\infty}^{\infty} \frac{e^{imx} dx}{(a^2 + x^2)^2} + 0. \quad (3.80)$$

As in Example 4, the imaginary part is zero because the integrand is an odd function leaving

$$\int_{-\infty}^{\infty} \frac{\cos mx \, dx}{(a^2 + x^2)^2} = \frac{\pi e^{-ma}}{2a^3} (1 + ma). \quad (3.81)$$

### 3.6 Integrals around the unit circle

We consider here integrals of the type  $\int_0^{2\pi} f(\cos \theta, \sin \theta) d\theta$ . Let us do this by example:

$$I = \int_0^{2\pi} \frac{d\theta}{a + \cos \theta}, \quad a > 1. \quad (3.82)$$

Take  $C$  as the unit circle  $z = e^{i\theta}$ . Therefore  $dz = ie^{i\theta} d\theta$  and

$$\begin{aligned} I &= \oint_C \frac{dz}{iz(a + \frac{1}{2}(z + z^{-1}))} \\ &= -2i \oint_C \frac{dz}{z^2 + 2az + 1}. \end{aligned} \quad (3.83)$$

The next task is to determine the roots of  $z^2 + 2az + 1 = 0$ .

$$z^2 + 2az + 1 = (z - \alpha^+)(z - \alpha^-) \quad (3.84)$$

where  $\alpha^\pm = -a \pm \sqrt{a^2 - 1}$ . Note that when  $a > 1$ , while  $\alpha^+$  lies *within*  $C$ ,  $\alpha^-$  lies *without*. Therefore we exclude the pole at  $z = \alpha^-$  and compute the residue of the integrand at  $z = \alpha^+$ , which is

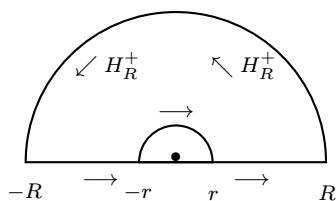
$$\frac{1}{\alpha^+ - \alpha^-}. \quad (3.85)$$

The Residue Theorem then gives

$$T = -2i \times 2\pi i \times \frac{1}{\alpha^+ - \alpha^-} = \frac{2\pi}{\sqrt{a^2 - 1}}. \quad (3.86)$$

### 3.7 Poles on the real axis

When an integrand has a pole on the real axis this means that it causes problems by sitting on the semicircular contour. Let us do this by example.



The contour is deformed by a small semi-circle of radius  $r$  centred at the origin that excludes the pole at  $z = 0$ . Following the direction of the arrows, the big semicircle of radius  $R$  is designated as  $H_R$  ( $\theta : 0 \rightarrow \pi$ ) and the little semicircle of radius  $r$  is designated as  $H_r$  ( $\theta : \pi \rightarrow 0$ ).

To calculate

$$\int_{-\infty}^{\infty} \frac{\sin x \, dx}{x} \quad (3.87)$$

we consider the complex integral

$$\oint_C \frac{e^{iz} \, dz}{z}. \quad (3.88)$$

The integrand has no poles in  $C$  because that  $z = 0$  is excluded in the above construction. Cauchy's Theorem can be invoked to give

$$0 = \oint_C \frac{e^{iz} \, dz}{z} = \int_{-R}^{-r} \frac{e^{ix} \, dx}{x} + \int_{H_r \leftarrow} \frac{e^{iz} \, dz}{z} + \int_r^R \frac{e^{ix} \, dx}{x} + \int_{H_R \rightarrow} \frac{e^{iz} \, dz}{z}. \quad (3.89)$$

Now we take the limit  $R \rightarrow \infty$  and, with  $m = 1$  Jordan's Lemma tell us that  $\int_{H_R} = 0$  because the only singularity is a pole and the integrand decays to zero as  $R \rightarrow \infty$ . Given that the small circle has the equation  $z = r(\cos \theta + i \sin \theta)$  for  $\theta : \pi \rightarrow 0$ , and noting that  $\sin \theta \geq 0$  in this range

$$\lim_{r \rightarrow 0} \int_{H_r} \frac{e^{iz} \, dz}{z} = i \lim_{r \rightarrow 0} \int_{\pi}^0 e^{-r \sin \theta} e^{ir \cos \theta} d\theta = -\pi i. \quad (3.90)$$

Taking the two limits  $R \rightarrow \infty$  and  $r \rightarrow 0$  together, we have

$$0 = \int_{-\infty}^{\infty} \frac{e^{ix}}{x} dx - \pi i + 0. \quad (3.91)$$

The real part of the integrand  $\cos x/x$  is odd so the contributions on  $(-\infty, 0)$  and  $(0, \infty)$  cancel leaving us with

$$\int_{-\infty}^{\infty} \frac{\sin x}{x} dx = \pi. \quad (3.92)$$

If the small circle is taken below the origin indented into the lower half-plane then its contribution is  $\pi i$  and, because the pole at  $z = 0$  is now included with residue unity, the contribution from the Residue Theorem is  $2\pi i$ . Thus we end up with the same result, as we should.